

Cleo Bench Problem 2

$$\text{Crazy } \int_0^\infty {}_3F_2\left(\begin{matrix} \frac{5}{8}, \frac{5}{8}, \frac{9}{8} \\ \frac{1}{2}, \frac{13}{8} \end{matrix} \middle| -x\right)^2 \frac{dx}{\sqrt{x}}$$

Derivation by Claude (Fable 5), closed-book*

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Problem

Evaluate in closed form

$$I = \int_0^\infty {}_3F_2\left(\begin{matrix} \frac{5}{8}, \frac{5}{8}, \frac{9}{8} \\ \frac{1}{2}, \frac{13}{8} \end{matrix} \middle| -x\right)^2 \frac{dx}{\sqrt{x}}.$$

Result

$$I = \frac{25 \Gamma\left(\frac{3}{4}\right)^2}{3\sqrt{\pi}} (2 \ln(1 + \sqrt{2}) - \sqrt{2}) = \frac{50 \pi^{3/2}}{3 \Gamma\left(\frac{1}{4}\right)^2} (2 \ln(1 + \sqrt{2}) - \sqrt{2})$$

$$I = 2.460685302449425108981967679044937181349580663151474781 \dots$$

(The two forms agree because $\Gamma(\frac{1}{4})\Gamma(\frac{3}{4}) = \pi\sqrt{2}$, hence $\Gamma(\frac{3}{4})^2 = 2\pi^2/\Gamma(\frac{1}{4})^2$. Also $2 \ln(1 + \sqrt{2}) = \ln(3 + 2\sqrt{2}) = 2 \operatorname{arcsinh} 1$.)

Derivation

Throughout put

$$a = \frac{5}{8}, \quad p = 2a = \frac{5}{4}.$$

The parameter pattern of the hypergeometric is exactly

$${}_3F_2\left(\begin{matrix} a, a, a + \frac{1}{2} \\ \frac{1}{2}, a + 1 \end{matrix} \middle| -x\right) =: F(x),$$

since $\frac{5}{8} = a$, $\frac{9}{8} = a + \frac{1}{2}$, $\frac{13}{8} = a + 1$. This pattern is what makes the problem tractable: the pair $(a, a + \frac{1}{2}; \frac{1}{2})$ triggers a quadratic (elementary) ${}_2F_1$ evaluation, and the pair $(a; a + 1)$ is a single Euler-type integration.

All powers of complex numbers below are **principal branches**; every complex number raised to a power will lie in the open right half-plane $\Re z > 0$ (or be a positive real), so no branch ambiguity ever arises.

*Problem originally posed on Mathematics Stack Exchange ([question 712798](#), CC BY-SA), famously answered by user Cleo. This derivation was produced independently, offline, without access to the published answer, as part of the Cleo benchmark.

Step 1. An elementary form for the ${}_2F_1$ kernel

Claim. For all real s ,

$${}_2F_1\left(a, a + \frac{1}{2}; \frac{1}{2}; -s^2\right) = \frac{(1 + is)^{-2a} + (1 - is)^{-2a}}{2} = \Re(1 + is)^{-2a}. \quad (1)$$

Proof. For $|z| < 1$ the binomial series gives

$$\frac{(1 - z)^{-2a} + (1 + z)^{-2a}}{2} = \sum_{m \text{ even}} \frac{(2a)_m}{m!} z^m = \sum_{k \geq 0} \frac{(2a)_{2k}}{(2k)!} z^{2k}.$$

Splitting the products over even/odd factors (the Pochhammer duplication identities)

$$(2a)_{2k} = 4^k (a)_k \left(a + \frac{1}{2}\right)_k, \quad (2k)! = 4^k k! \left(\frac{1}{2}\right)_k,$$

we get $\sum_k \frac{(a)_k (a + \frac{1}{2})_k}{(\frac{1}{2})_k k!} z^{2k} = {}_2F_1(a, a + \frac{1}{2}; \frac{1}{2}; z^2)$ for $|z| < 1$.

Now fix the branch discussion for $z = is$, $s \in \mathbb{R}$. The left side, as a function of s , is $\varphi(s^2)$ where $\varphi(w) = {}_2F_1(a, a + \frac{1}{2}; \frac{1}{2}; -w)$ is analytic on $\mathbb{C} \setminus (-\infty, -1]$; since $-s^2 \leq 0$ we stay in its domain, so $s \mapsto \varphi(s^2)$ is real-analytic on all of \mathbb{R} . The right side of (1) is real-analytic on \mathbb{R} as well, because $1 \pm is$ lies in the closed right half-plane minus 0 for every real s , away from the cut of the principal power. Both sides agree for $|s| < 1$ by the series computation, hence for all real s by analytic continuation. Finally $(1 + is)^{-2a} = (1 - is)^{-2a}$ for real s , so the average is $\Re(1 + is)^{-2a}$. \square

Step 2. Euler integral for the lower-parameter shift $a \mapsto a + 1$

Since $\frac{(a)_n}{(a+1)_n} = \frac{a}{a+n} = a \int_0^1 t^{a+n-1} dt$, termwise integration of the (uniformly convergent) series for $0 \leq x < 1$ gives

$$F(x) = a \int_0^1 t^{a-1} {}_2F_1\left(a, a + \frac{1}{2}; \frac{1}{2}; -xt\right) dt.$$

Both sides are analytic in x on $\Omega = \mathbb{C} \setminus (-\infty, -1]$ — the left side is the principal branch of the ${}_3F_2$; the right side is analytic on Ω because for $x \in \Omega$ and $t \in (0, 1]$ we have $xt \in \Omega$ (the product xt can only land on $(-\infty, -1]$ if x itself is real ≤ -1), the integrand is jointly continuous and locally uniformly bounded, so Morera/dominated convergence applies. Agreement on $[0, 1)$ therefore extends to all $x \geq 0$.

Inserting (1) with $s = \sqrt{xt}$ and substituting $t = v^2$ ($dt = 2v dv$):

$$F(x) = 2a \int_0^1 v^{2a-1} \Re(1 + iv\sqrt{x})^{-2a} dv, \quad x \geq 0. \quad (2)$$

(Representation (2) was verified numerically to 50 digits at $x = 0.7, 3.3, 100$; see the verification section.)

Step 3. Square, substitute $x = y^2$, and interchange (Fubini)

With $x = y^2$, $dx/\sqrt{x} = 2 dy$, and writing the square of (2) as a double integral,

$$I = 2 \int_0^\infty F(y^2)^2 dy = 2(2a)^2 \int_0^\infty \iint_{[0,1]^2} (uv)^{2a-1} \Re(1 + iuy)^{-p} \Re(1 + ivy)^{-p} du dv dy.$$

To swap the y -integral with the (u, v) -integral we check absolute integrability. Since $|\Re(1 + iuy)^{-p}| \leq |1 + iuy|^{-p} = (1 + u^2y^2)^{-p/2} \leq 1$, with $m = \max(u, v)$,

$$\int_0^\infty |\Re(1 + iuy)^{-p} \Re(1 + ivy)^{-p}| dy \leq \int_0^\infty (1 + m^2y^2)^{-p/2} dy = \frac{c_1}{m}, \quad c_1 = \int_0^\infty \frac{ds}{(1 + s^2)^{5/8}} < \infty$$

(c_1 converges because $2 \cdot \frac{5}{8} = \frac{5}{4} > 1$), and

$$\iint_{[0,1]^2} (uv)^{1/4} \frac{c_1}{\max(u, v)} du dv = 2c_1 \int_0^1 u^{-3/4} \left(\int_0^u v^{1/4} dv \right) du = \frac{16c_1}{15} < \infty.$$

(This bound also proves $I < \infty$.) Hence, by Fubini,

$$I = 2(2a)^2 \iint_{[0,1]^2} (uv)^{2a-1} W(u, v) du dv, \quad W(u, v) = \int_0^\infty \Re(1 + iuy)^{-p} \Re(1 + ivy)^{-p} dy. \quad (3)$$

Step 4. The kernel integral $W(u, v)$ in closed form

Write $A = (1 + iuy)^{-p}$, $B = (1 + ivy)^{-p}$; for real y , $\bar{B} = (1 - ivy)^{-p}$, and

$$\Re A \Re B = \frac{1}{2} \Re(AB) + \frac{1}{2} \Re(A\bar{B}) \implies W = \frac{1}{2} \Re J_+ + \frac{1}{2} \Re J_-,$$

$$J_+ = \int_0^\infty (1 + iuy)^{-p} (1 + ivy)^{-p} dy, \quad J_- = \int_0^\infty (1 + iuy)^{-p} (1 - ivy)^{-p} dy.$$

(a) $\Re J_+ = 0$. For $\kappa > 0$ the principal power $(1 + i\kappa y)^{-p}$ is analytic in y except on the cut $\{y = it : t \geq 1/\kappa\}$ (where $1 + i\kappa y \in (-\infty, 0]$), which lies on the **positive** imaginary axis. So the integrand of J_+ is analytic on the closed fourth quadrant. On the arc $|y| = R$, $\arg y \in [-\frac{\pi}{2}, 0]$, the integrand is $O(R^{-2p}) = O(R^{-5/2})$, so the arc contributes $O(R^{-3/2}) \rightarrow 0$. By Cauchy's theorem we may rotate the ray of integration to the negative imaginary axis, $y = -i\tau$:

$$J_+ = -i \int_0^\infty (1 + u\tau)^{-p} (1 + v\tau)^{-p} d\tau \in i\mathbb{R} \implies \Re J_+ = 0.$$

(b) $\Re J_-$. For $y \geq 0$ both $1 + iuy$ and $1 - ivy$ lie in the open right half-plane, where principal powers multiply: $(z_1 z_2)^{-p} = z_1^{-p} z_2^{-p}$ whenever $|\arg z_1|, |\arg z_2| < \frac{\pi}{2}$. Completing the square,

$$(1 + iuy)(1 - ivy) = 1 + uv y^2 + i(u - v)y = uv \left[(y + i\delta)^2 + \sigma^2 \right], \quad \delta = \frac{u - v}{2uv}, \quad \sigma = \frac{u + v}{2uv},$$

which is consistent at $y = 0$ because $uv(\sigma^2 - \delta^2) = \frac{(u+v)^2 - (u-v)^2}{4uv} = 1$; note $\sigma > |\delta|$. Since $\Re[(y + i\delta)^2 + \sigma^2] = \frac{1 + uv y^2}{uv} > 0$ and $uv > 0$, principal branches again factor, and

$$J_- = (uv)^{-p} \int_\Gamma (w^2 + \sigma^2)^{-p} dw,$$

where Γ is the horizontal ray from $i\delta$ to $i\delta + \infty$ ($w = y + i\delta$). The principal power $(w^2 + \sigma^2)^{-p}$ is analytic off $\{w = it : |t| \geq \sigma\}$, and the strip $|\Im w| \leq |\delta| < \sigma$, $\Re w \geq 0$, avoids that set. Deform Γ to the vertical segment from $i\delta$ down to 0 followed by $[0, \infty)$; the far vertical connector at $\Re w = R$ contributes $O(R^{-5/2}) \cdot |\delta| \rightarrow 0$. On the segment $w = i\tau$ (τ between 0 and δ) we have $w^2 + \sigma^2 = \sigma^2 - \tau^2 \in (0, \sigma^2]$, so the integrand is **real** while $dw = i d\tau$ — that piece is purely imaginary. Hence

$$\Re J_- = (uv)^{-p} \int_0^\infty (y^2 + \sigma^2)^{-p} dy = (uv)^{-p} \sigma^{1-2p} \frac{\sqrt{\pi} \Gamma\left(p - \frac{1}{2}\right)}{2\Gamma(p)} = \frac{\sqrt{\pi} \Gamma\left(p - \frac{1}{2}\right)}{2\Gamma(p)} \frac{2^{2p-1} (uv)^{p-1}}{(u+v)^{2p-1}},$$

using the Beta evaluation $\int_0^\infty (y^2 + \sigma^2)^{-p} dy = \frac{1}{2} \sigma^{1-2p} B(\frac{1}{2}, p - \frac{1}{2})$ (substitute $y = \sigma \tan \theta$), valid since $p = \frac{5}{4} > \frac{1}{2}$.

So

$$W(u, v) = \frac{\sqrt{\pi} \Gamma(p - \frac{1}{2})}{4\Gamma(p)} \frac{2^{2p-1} (uv)^{p-1}}{(u+v)^{2p-1}}. \quad (4)$$

(Identity (4) was verified numerically to ~ 50 digits at several (u, v) ; see below.)

Step 5. Assembling: an elementary double integral

Insert (4) into (3). With $a = \frac{5}{8}$, $p = \frac{5}{4}$ we have $(uv)^{2a-1}(uv)^{p-1} = (uv)^{1/2}$ and $2p - 1 = \frac{3}{2}$:

$$I = (2a)^2 \frac{\sqrt{\pi} \Gamma(\frac{3}{4})}{2\Gamma(\frac{5}{4})} 2^{3/2} D, \quad D = \iint_{[0,1]^2} \frac{\sqrt{uv}}{(u+v)^{3/2}} du dv.$$

By the $u \leftrightarrow v$ symmetry and the scaling $v = ut$ on the region $v < u$,

$$D = 2 \int_0^1 \int_0^u \frac{\sqrt{uv}}{(u+v)^{3/2}} dv du = 2 \int_0^1 u^{1/2} du \int_0^1 \frac{\sqrt{t}}{(1+t)^{3/2}} dt = \frac{4}{3} \int_0^1 \frac{\sqrt{t}}{(1+t)^{3/2}} dt.$$

Substituting $t = \sinh^2 w$ turns the last integral into $2 \int \tanh^2 w dw = 2(w - \tanh w)$, so

$$\int_0^1 \frac{\sqrt{t}}{(1+t)^{3/2}} dt = 2 \left(\operatorname{arcsinh} 1 - \frac{1}{\sqrt{2}} \right) = 2 \ln(1 + \sqrt{2}) - \sqrt{2}, \quad D = \frac{4}{3} \left(2 \ln(1 + \sqrt{2}) - \sqrt{2} \right).$$

Step 6. Final constants

$$I = \frac{25}{16} \cdot \frac{\sqrt{\pi} \Gamma(\frac{3}{4})}{2\Gamma(\frac{5}{4})} \cdot 2^{3/2} \cdot \frac{4}{3} \left(2 \ln(1 + \sqrt{2}) - \sqrt{2} \right) = \frac{25 \cdot 2^{7/2}}{96} \cdot \frac{\sqrt{\pi} \Gamma(\frac{3}{4})}{\Gamma(\frac{5}{4})} \left(2 \ln(1 + \sqrt{2}) - \sqrt{2} \right).$$

Now $\Gamma(\frac{5}{4}) = \frac{1}{4} \Gamma(\frac{1}{4})$ and the reflection formula $\Gamma(\frac{1}{4}) \Gamma(\frac{3}{4}) = \frac{\pi}{\sin(\pi/4)} = \pi\sqrt{2}$ give $\frac{\sqrt{\pi} \Gamma(3/4)}{\Gamma(5/4)} = \frac{4\sqrt{2} \pi^{3/2}}{\Gamma(1/4)^2}$, hence

$$I = \frac{25 \cdot 2^{7/2} \cdot 4\sqrt{2}}{96} \cdot \frac{\pi^{3/2}}{\Gamma(\frac{1}{4})^2} \left(2 \ln(1 + \sqrt{2}) - \sqrt{2} \right) = \frac{50 \pi^{3/2}}{3 \Gamma(\frac{1}{4})^2} \left(2 \ln(1 + \sqrt{2}) - \sqrt{2} \right),$$

and equivalently, using $\Gamma(\frac{3}{4})^2 = 2\pi^2/\Gamma(\frac{1}{4})^2$,

$$I = \frac{25 \Gamma(\frac{3}{4})^2}{3\sqrt{\pi}} \left(2 \ln(1 + \sqrt{2}) - \sqrt{2} \right).$$

■

Remark. The same computation evaluates the whole family $\int_0^\infty {}_3F_2(a, a, a + \frac{1}{2}; \frac{1}{2}, a + 1; -x)^2 x^{-1/2} dx$ for $\frac{1}{2} < a < 1$ (so that $p = 2a$ satisfies $p > \frac{1}{2}$ for Step 4 and the Fubini bounds hold); the double integral $\iint (uv)^{3a-2} (u+v)^{1-4a} du dv$ is again elementary-hypergeometric. Wait — for general a the exponents are $(uv)^{2a-1+p-1} = (uv)^{3a-2}$ and $(u+v)^{4a-1}$; the special value $a = \frac{5}{8}$ makes it the pleasant $\iint \sqrt{uv} (u+v)^{-3/2}$. (*This remark is illustrative only; the boxed result for $a = 5/8$ is what was proved above.*)

Numerical verification

All computations with `mpmath`. Working precision `mp.dps = 60` (spot checks at 50).

1. **Representation (2)** vs `mp.hyp3f2(5/8,5/8,9/8,1/2,13/8,-x)` at $x = 0.7, 3.3, 100$: differences 0, 0, 8×10^{-53} at 50 dps.
2. **Kernel identity (4)** vs direct quadrature of $W(u, v) = \int_0^\infty \Re(1 + iuy)^{-5/4} \Re(1 + ivy)^{-5/4} dy$ at $(u, v) = (0.3, 0.9), (0.5, 0.5), (0.05, 0.77)$: differences $\lesssim 3 \times 10^{-51}$.
3. **Direct evaluation of I** at 60 dps using `mp.hyp3f2` itself (no ingredient of the derivation): $I = \int_0^1 F(x)^2 x^{-1/2} dx + \int_0^1 F(1/w)^2 w^{-3/2} dw$, computed as $2 \int_0^1 F(y^2)^2 dy$ (smooth) plus, with $w = z^4$, $4 \int_0^1 F(z^{-4})^2 z^{-3} dz$ (integrand $\sim z^2 \log^2 z$ at 0), each by `tanh-sinh` quadrature (`maxdegree=9`):

$$I_{\text{num}} = 2.460685302449425108981967679044937181349580663151474781 \dots$$

4. **Closed form** at 60 dps:

$$\frac{25 \Gamma(3/4)^2}{3\sqrt{\pi}} (2 \ln(1 + \sqrt{2}) - \sqrt{2})$$

$$= 2.460685302449425108981967679044937181349580663151474781 \dots$$

Relative difference $|I_{\text{num}} - I_{\text{cf}}|/I_{\text{cf}} \approx 1.3 \times 10^{-61}$: **60 significant digits agree** (the full working precision).

5. **Independent cross-check (Mellin–Parseval)**. With $\tilde{F}(s) = \frac{\Gamma(s)\Gamma(\frac{5}{8}-s)^2\Gamma(\frac{9}{8}-s)\Gamma(\frac{1}{2})\Gamma(\frac{13}{8})}{\Gamma(\frac{5}{8})^2\Gamma(\frac{9}{8})\Gamma(\frac{1}{2}-s)\Gamma(\frac{13}{8}-s)}$ (the Mellin transform of F), Parseval gives $I = \frac{1}{\pi} \int_0^\infty \Re[\tilde{F}(\frac{1}{4} + it)\tilde{F}(\frac{1}{4} - it)] dt$; after the cancellation of $\Gamma(s)$ and $\Gamma(\frac{1}{2} - s)$ in the product this Gamma-integral was evaluated numerically at 60 dps and agrees with the closed form to 6×10^{-61} . This route uses none of Steps 1–6.

Scripts: `step1_checks.py`, `step2_direct.py`, `step3_highprec.py`, `step4_digits.py` in the scratch work directory.

Notes

- The proof is self-contained: the only inputs are the binomial series, the Pochhammer duplication identities, Cauchy’s theorem (two explicitly justified contour deformations, with the decay rate $|y|^{-5/2}$ making all arcs vanish), the Beta integral $\int_0^\infty (y^2 + \sigma^2)^{-p} dy$, and one Fubini interchange with an explicit absolutely-convergent majorant (which simultaneously proves $I < \infty$).
- Branch handling: every principal power occurs at an argument in the closed right half-plane off the cut, and both factorization steps $(z_1 z_2)^{-p} = z_1^{-p} z_2^{-p}$ are applied only when $|\arg z_i| < \pi/2$, where they are valid.
- The analytic-continuation arguments in Steps 1–2 (from $|x| < 1$ to $x \geq 0$) are standard and fully justified (domains are connected, both sides analytic, agreement on a set with a limit point); they were additionally confirmed numerically.

- The Mellin–Parseval computation in the verification section is only a cross-check, not part of the proof; its own hypotheses (strip of analyticity $0 < \Re s < \frac{5}{8}$ containing $\Re s = \frac{1}{4}$, L^2/L^1 conditions from the $e^{-2\pi|t|}$ decay of the Gamma product) do hold, but the derivation does not rely on it.
- No step is conjectural; I consider the result proved.